



# Current CaCO<sub>3</sub> dissolution at the seafloor caused by anthropogenic CO<sub>2</sub>

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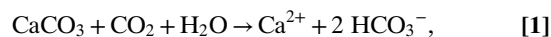
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**Oceanic uptake of anthropogenic CO<sub>2</sub> leads to decreased pH, carbonate ion concentration, and saturation state with respect to CaCO<sub>3</sub> minerals, causing increased dissolution of these minerals at the deep seafloor. This additional dissolution will figure prominently in the neutralization of man-made CO<sub>2</sub>. However, there has been no concerted assessment of the current extent of anthropogenic CaCO<sub>3</sub> dissolution at the deep seafloor. Here, recent databases of bottom-water chemistry, benthic currents, and CaCO<sub>3</sub> content of deep-sea sediments are combined with a rate model to derive the global distribution of benthic calcite dissolution rates and obtain primary confirmation of an anthropogenic component. By comparing preindustrial with present-day rates, we determine that significant anthropogenic dissolution now occurs in the western North Atlantic, amounting to 40–100% of the total seafloor dissolution at its most intense locations. At these locations, the calcite compensation depth has risen ~300 m. Increased benthic dissolution was also revealed at various hot spots in the southern extent of the Atlantic, Indian, and Pacific Oceans. Our findings place constraints on future predictions of ocean acidification, are consequential to the fate of benthic calcifiers, and indicate that a by-product of human activities is currently altering the geological record of the deep sea.**

ocean acidification | seafloor | CaCO<sub>3</sub> | dissolution | anthropogenic CO<sub>2</sub>

**S**eafloor dissolution of CaCO<sub>3</sub> minerals will constitute a primary feedback to ocean acidification over timescales of centuries to tens of millennia (1). The overall dissolution reaction is as follows:



where CaCO<sub>3</sub> denotes solid carbonate in bottom sediments, mainly as calcite. This process is termed geochemical carbonate compensation. Consequently, CO<sub>2</sub> entering the ocean, including that of anthropogenic origin, can be neutralized permanently by conversion to dissolved bicarbonate ions (HCO<sub>3</sub><sup>-</sup>).

The upper oceans are wholly supersaturated with respect to calcite, despite the current acidification. Largely because of the increasing solubility of calcite with pressure, the deeper oceans become undersaturated, whereafter the rate of reaction 1 increases with oceanographic depth. The depth where undersaturation first occurs is the calcite saturation depth (CSD) (2). The preindustrial oceans contained CO<sub>2</sub> acquired from the atmosphere, from marine volcanism, and from oxic organic matter decay. At the same time, calcifying organisms precipitated CaCO<sub>3</sub> shells that settled to the seafloor upon their death. Combined, these processes and reaction 1 lead to decreasing CaCO<sub>3</sub> content in sediments below the CSD. The depth where the deposition rate of CaCO<sub>3</sub> is exactly balanced by reaction 1 is called the calcite compensation depth (CCD) (2), although operationally it is commonly defined as the depth where the

sediment CaCO<sub>3</sub> content falls below 10%. The snowline (2) denotes the depth below which no CaCO<sub>3</sub> is found in the sediment; the snowline and CCD coincide at steady state (2).

The geological record contains numerous examples of deep-sea CaCO<sub>3</sub> dissolution events driven by natural acidification, for example, at ~56 My BP, known as the Paleocene–Eocene Thermal Maximum (3–5). During these events, CaCO<sub>3</sub> disappeared from deep-sea sediments, where it had previously accumulated, burial rates dropped, and the snowline shoaled. The same scenario has been predicted for the Anthropocene oceans (6, 7), but no estimates of increased deep-sea sediment CaCO<sub>3</sub> dissolution have been published. This lack of documentation might be attributed to the restricted penetration of anthropogenic CO<sub>2</sub> to shallow depths, but this premise is contradicted by observed changes in the carbonate chemistry of the deeper oceans (8–10) and as reported here. Both numerical ocean models and the presence in abyssal waters of transient tracers produced almost entirely after the end of the 1940s, that is, chlorofluorocarbons and polychlorinated biphenyls, strongly imply the presence of anthropogenic CO<sub>2</sub> in the deep oceans. Anthropogenic CO<sub>2</sub> in deep and bottom waters is simply very difficult to measure because it is a small signal superimposed on a large natural background concentration. Early changes in

## Significance

The geological record contains numerous examples of “greenhouse periods” and ocean acidification episodes, where the spreading of corrosive (CO<sub>2</sub>-enriched) bottom waters enhances the dissolution of CaCO<sub>3</sub> minerals delivered to the seafloor or contained within deep-sea sediments. The dissolution of sedimentary CaCO<sub>3</sub> neutralizes excess CO<sub>2</sub>, thus preventing runaway acidification, and acts as a negative-feedback mechanism in regulating atmospheric CO<sub>2</sub> levels over timescales of centuries to millennia. We report an observation-based indication and quantification of significant CaCO<sub>3</sub> dissolution at the seafloor caused by man-made CO<sub>2</sub>. This dissolution is already occurring at various locations in the deep ocean, particularly in the northern Atlantic and near the Southern Ocean, where the bottom waters are young and rich in anthropogenic CO<sub>2</sub>.

Author contributions: O.S., B.P.B., and A.M. designed research; O.S. performed research; O.S., B.P.B., A.M., C.J., D.S.T., B.K.A., and R.M.K. contributed new reagents/analytic tools; O.S., B.P.B., C.J., D.S.T., B.K.A., and R.M.K. analyzed data; and O.S., B.P.B., A.M., C.J., D.S.T., B.K.A., and R.M.K. wrote the paper.

The authors declare no conflict of interest.

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This article contains supporting information online at [www.pnas.org/lookup/suppl/doi:10.1073/pnas.1804250115/-DCSupplemental](http://www.pnas.org/lookup/suppl/doi:10.1073/pnas.1804250115/-DCSupplemental).

CaCO<sub>3</sub> content of sediments caused by anthropogenic acidification are also extremely difficult to detect through changes in either solid mass or dissolution indices.

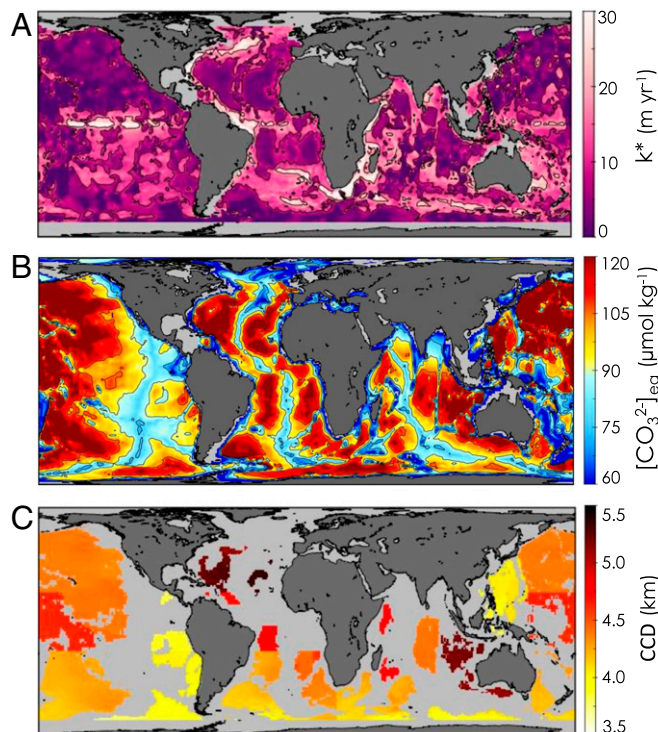
Given that sediment monitoring is unlikely to yield immediate evidence of anthropogenic calcite dissolution at the seafloor, we employ a different approach and compare the rate of dissolution at the seafloor under preindustrial (~1800 AD) and modern (2002 AD) benthic conditions. In accord with recent work (11, 12) on the dissolution kinetics of calcite beds, the dissolution rate ( $r$ ) at any depth between the CSD and the CCD is given by the following (11):

$$r = k^* \left( [\text{CO}_3^{2-}]_{\text{eq}} - [\text{CO}_3^{2-}]_{\text{sw}} \right), \quad [2]$$

where  $[\text{CO}_3^{2-}]_{\text{eq}}$  is the calcite-equilibrium carbonate ion concentration (Fig. 1B),  $[\text{CO}_3^{2-}]_{\text{sw}}$  is the carbonate ion concentration in bottom waters, and  $k^*$  is the overall CO<sub>3</sub><sup>2-</sup> mass transfer coefficient (11):

$$k^* = k_s \beta (k_s + \beta)^{-1}, \quad [3]$$

where  $k_s$  is the sediment-side mass transfer coefficient, which characterizes dissolution and transport (diffusion) of the carbonate ion within the sediment, and  $\beta$  is the water-side mass transfer coefficient (13). In this formulation,  $k^*$  tends toward the value of the smallest rate-limiting mass transfer coefficient,  $\beta$  or  $k_s$ , without ever reaching it, as predicted by theory and validated by observations. For further explanation with regard to the derivation of Eqs. 2 and 3, please see *SI Appendix*.



**Fig. 1.** Chemical and physical parameters for calcite dissolution. (A) Overall mass transfer coefficient  $k^*$  for CaCO<sub>3</sub> dissolution, (B) current bottom-water saturation concentration  $[\text{CO}_3^{2-}]_{\text{eq}}$  at in situ temperature, salinity, and pressure, and (C) present-day calcite compensation depth (CCD), regionally averaged as described in *Materials and Methods*, where the light gray identifies areas where the CCD is deeper than the depth of the seafloor, that is, where calcite can undergo net sediment accumulation. The CCD is not computed in the Arctic and Southern Oceans (south of 60°S).

With any Anthropocene (14) rise in the CCD,  $r$  is supplemented by dissolution of previously deposited CaCO<sub>3</sub> between the new and old CCD positions (*Materials and Methods*). Hence, the difference between the calculated  $r$  over the deep seafloor for preindustrial and current conditions reflects the impacts of anthropogenic acidification. To perform this calculation globally, we require the spatial distributions of  $k^*$ ,  $[\text{CO}_3^{2-}]_{\text{eq}}$ ,  $[\text{CO}_3^{2-}]_{\text{sw}}$ , and the CaCO<sub>3</sub> contents of surface sediments, whose distribution is available from a new database (15, 16) and displayed in *SI Appendix*, Fig. S24.

The distribution of  $k^*$  at the seafloor is shown in Fig. 14. The water-side mass transfer coefficient,  $\beta$ , appears in Eq. 3 because the seafloor is covered by a water layer through which solute transport occurs via molecular diffusion (13), termed the diffusive boundary layer (DBL). The existence of the DBL has been amply illustrated by previous research (17–19).  $\beta$  is calculated as the ratio of the diffusion coefficient of CO<sub>3</sub><sup>2-</sup> ( $D_{\text{CO}_3^{2-}}$ ), at in situ temperature ( $T$ ) and salinity ( $S_P$ ), to the thickness of the DBL ( $Z_{\text{DBL}}$ ), as  $\beta$  is effectively independent of pressure (*Materials and Methods*). Whereas  $D_{\text{CO}_3^{2-}}$  does not vary widely,  $Z_{\text{DBL}}$  depends on the flow velocities at the seafloor (13), but an ocean-wide distribution of  $Z_{\text{DBL}}$  has never been reported in the literature. Herein, we use a global bottom-current velocity ( $U$ ) model to derive in situ  $\beta$  values (*Materials and Methods*). *SI Appendix*, Fig. S3C provides an ocean-wide distribution of  $\beta$  and shows that this parameter is high on the east side of continents, beneath the Equator, and on the northern fringe of the Southern Ocean, all areas with enhanced bottom currents. Conversely,  $\beta$  is very small over much of the abyssal ocean due to relatively sluggish flow. The magnitude of  $k_s$  is derived from available experiments and calculated as a function of the CaCO<sub>3</sub> content in surface sediments (*Materials and Methods*). The distribution of  $k_s$  at the seafloor is shown in *SI Appendix*, Fig. S2B. Overall, the rate of deep-sea CaCO<sub>3</sub> dissolution,  $r$ , is largely controlled by  $\beta$ , rather than  $k_s$  (11, 12), except in regions of high bottom currents or where sediments are CaCO<sub>3</sub>-poor, such as the North Pacific or the Southern Ocean (*SI Appendix*, Fig. S5).

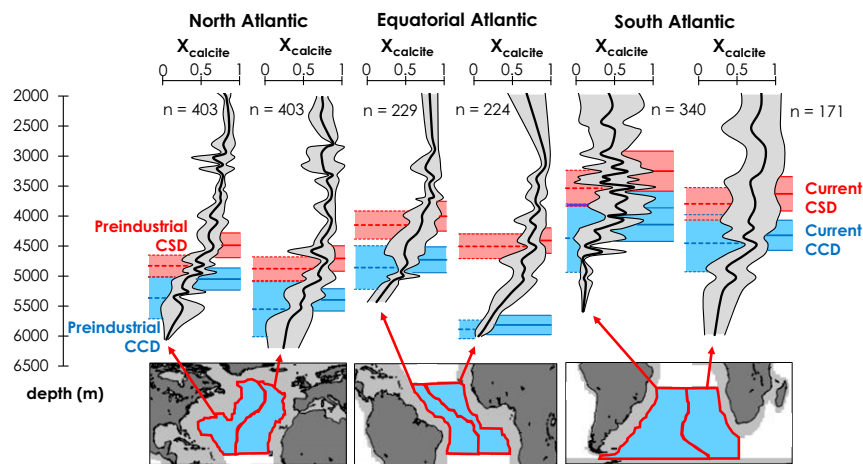
The benthic distribution of  $[\text{CO}_3^{2-}]_{\text{eq}}$  in Eq. 2 is calculated by dividing the stoichiometric solubility constant of calcite,  $K_{sp}^*$ , at in situ conditions, by the Ca<sup>2+</sup> concentration of the oceans. Fig. 1B illustrates the resulting map of  $[\text{CO}_3^{2-}]_{\text{eq}}$ .  $[\text{CO}_3^{2-}]_{\text{eq}}$  is commonly near 75 μmol·kg<sup>-1</sup> on top of oceanic ridges (~2,500 m) but increases to values as high as ~140 μmol·kg<sup>-1</sup> on the abyssal plains (~6,000 m).

Next, we need to estimate the position of the preindustrial CCD to calculate the amount of dissolution below this depth. As stated earlier, the CCD and snowline are coincident at oceanic steady state. In addition, there is no evidence that the calcite snowline has, as yet, migrated measurably due to anthropogenic dissolution (7, 20). Therefore, we set the preindustrial CCD to the current snowline depth that we estimate from the CaCO<sub>3</sub> contents of sediments for each basin (*Materials and Methods*). With these local preindustrial CCD values, the flux ( $F$ ) of CaCO<sub>3</sub> at that depth, and corresponding grid point, can be estimated (2), and the dissolution rate below the CCD, set to the value of  $F$ , mapped.

We note that calcifying organisms exhibit various responses to elevated pCO<sub>2</sub> conditions, due to the influence of other climate change-related effects, such as the warming of waters that counteracts acidification (21–24). In the absence of unequivocal evidence of an immediate increase or decrease in the total calcification rate in the pelagic oceans, we assume that the flux of calcitic particles reaching the CCD ( $F$ ) has remained invariant since the preindustrial era. This is supported by what took place in similar acidification events in the geological record, such as the Paleocene–Eocene Thermal Maximum, where no obvious plankton carbonate productivity reduction is readily detectable (25). Thus, the “current” CCD was computed from the fixed (known)  $F$  and the present-day  $[\text{CO}_3^{2-}]_{\text{sw}}$ , and is represented in Fig. 1C.

Finally, the preindustrial distribution of bottom-water  $[\text{CO}_3^{2-}]_{\text{sw}}$  (Fig. 24) is computed using estimated preindustrial, deep-ocean





**Fig. 4.** Atlantic Ocean sediment calcite content profiles and calcite marker horizons. Calcite fraction in dry sediments ( $X_{\text{calcite}}$ ),  $\pm 1$  SD, for the basins defined by the maps below the plots. The number of measurements ( $n$ ) comprised within each basin is also reported. The preindustrial calcite compensation depth (CCD) and saturation depth (CSD) are represented with dashed horizontal lines (on the *Left* of each profile), whereas the solid horizontal lines correspond to current values (on the *Right* of each profile). Both the CCD and CSD are reported along with their uncertainties.

## Discussion

On average, we find that  $F$ , the flux of  $\text{CaCO}_3$  that must reach the CCD to explain the observed  $\text{CaCO}_3$  contents in surficial sediments (*Materials and Methods*) is  $0.17 \pm 0.07 \text{ mol}\cdot\text{m}^{-2}\cdot\text{y}^{-1}$ . This value is equivalent to  $\text{CaCO}_3$  fluxes to the deep ocean derived from previous studies based on excess total alkalinity ( $\text{TA}^*$ ) budgets (29, 30). Integrating over the surface of the seafloor, we find a global calcite downward flux at the CCD of  $54 \pm 17 \times 10^{12} \text{ mol}\cdot\text{y}^{-1}$  (i.e.,  $0.7 \pm 0.2 \text{ Gt C}\cdot\text{y}^{-1}$ ), and a current global seafloor  $\text{CaCO}_3$  dissolution flux of  $32 \pm 12 \times 10^{12} \text{ mol}\cdot\text{y}^{-1}$  of  $\text{CaCO}_3$  (i.e.,  $0.4 \pm 0.1 \text{ Gt C}\cdot\text{y}^{-1}$ ). Both of these values are similar to and more precise than the observation-based estimates presented in ref. 29. Our global dissolution estimate represents only  $22 \pm 8\%$  of the total estimated  $\text{CaCO}_3$  dissolution in the marine environment each year, that is,  $144 \times 10^{12} \text{ mol}\cdot\text{y}^{-1}$ , or  $1.7 \text{ Gt C}\cdot\text{y}^{-1}$  (31). The remainder of the dissolution occurs primarily in the water column and, possibly, after burial through metabolic dissolution or later diagenesis at shallow to mid depths.

As stated and justified earlier, we assumed that the calcite rain rate has remained invariant since preindustrial times. Although a decrease in the precipitation rate of calcite might be expected under more acidic conditions (32), thus reducing the calcite export from the surface ocean and the rain rate to deep-sea sediments, no such trend is clearly observed today (33). Testing the sensitivity of the computed sediment dissolution rates ( $r$ ) to a reduction in the calcite downward flux ( $F$ ), we find that when  $F$  is decreased by 10%, the global dissolution rate increases by 2.5%. Likewise, when  $F$  is decreased by 20%, the global dissolution rate increases by 3.1%. Thus, counterintuitively, if less calcite is currently delivered to the seafloor than during the preindustrial era, more calcite is being dissolved at the SWI because a weaker  $F$  causes the CCD to respond faster to changes in  $[\text{CO}_3^{2-}]_{\text{sw}}$ , as explained in ref. 2, and in the case of the present acidification, to rise faster. Sediments that were above the preindustrial CCD and are now below the current CCD are subject to the greatest dissolution rate increase, as can be seen when comparing Eqs. 12–14 in *Materials and Methods*. In other words, if  $F$  has remained constant,  $\sim 2\%$  of the seafloor is currently comprised between the preindustrial and current CCD, but this fraction increases to 9.3% when  $F$  is reduced by 20% relative to its preindustrial value.

Our results provide a tangible indication that significant anthropogenic calcite dissolution is currently occurring at the seafloor. This dissolution will eventually change the  $\text{CaCO}_3$  accumulation patterns and rates in the oceans, while mitigating runaway ocean acidification. Any future model of oceanic acidification needs to reproduce our observations to assure its validity. The consequences of this anthropogenic dissolution to the ecology of benthic calcifiers has not been determined with certainty, but could be as substantial as it is for pelagic calcifiers (34). Chemical burndown of previously

deposited carbonate-rich sediments has already begun and will intensify and spread over vast areas of the seafloor during the next decades and centuries, thus altering the geological record of the deep sea. The deep-sea benthic environment, which covers  $\sim 60\%$  of our planet, has indeed entered the Anthropocene.

## Materials and Methods

**Bottom-Water Chemistry.**  $[\text{CO}_3^{2-}]_{\text{sw}}$  values were computed using the Matlab version of the CO2SYS program (35, 36), based on GLODAPv2.2016b  $1^\circ \times 1^\circ$  global mapped climatologies (13, 37) at the deepest resolved layers of the GLODAPv2.2016b model (shown in *SI Appendix, Fig. S1*), assumed to be near the seafloor. The GLODAPv2.2016b variables used for our computation are represented in *SI Appendix, Fig. S1* and include in situ bottom-water TA, DIC,  $T$ ,  $S_p$ , as well as soluble reactive phosphate ([SRP]) and dissolved inorganic silica ([DSi]) concentrations. Current  $[\text{CO}_3^{2-}]_{\text{sw}}$  were derived using GLODAPv2.2016b bottom-water DIC data normalized to the year 2002 (*SI Appendix, Fig. S1G*), whereas preindustrial  $[\text{CO}_3^{2-}]_{\text{sw}}$  values were computed using the GLODAPv2.2016b bottom-water preindustrial DIC data (13) (*SI Appendix, Fig. S1H*). For computational purposes, we assume that bottom-water TA has not changed since the end of the preindustrial period ( $\sim 1800$  AD), in accordance with recent studies that predict an insignificant global mean deep-water TA anomaly ( $\Delta\text{TA} \sim 2.5 \mu\text{mol}\cdot\text{kg}^{-1}$  relative to the mean concentration from 1950 to 1959) before  $\sim 2060$  AD, although some changes might already be detectable (during the 2010–2020 decade) in the deep North Atlantic (38). We computed the TA flux generated by anthropogenic dissolution of  $\text{CaCO}_3$  sediments since the preindustrial times. Assuming that this TA mixes in the water column up to the permanent thermocline, we found that it has resulted in a world-averaged TA increase of less than  $0.1 \mu\text{mol}\cdot\text{kg}^{-1}$ , and of  $\sim 3 \mu\text{mol}\cdot\text{kg}^{-1}$  above the hot spot of anthropogenic dissolution in the Northwest Atlantic; both values are at or below the current analytical uncertainty.  $[\text{Ca}^{2+}]_{\text{eq}}$  was estimated from the equations presented in ref. 2, using  $[\text{Ca}^{2+}]$  computed from  $S_p$  following ref. 39, the stoichiometric solubility product of calcite ( $K_{sp}^*$ ) (40) at in situ  $T$  and  $S_p$ , and the pressure ( $P$ ) derived from GEBCO bathymetry (41).

**$\text{CaCO}_3$  Distribution in Deep-Sea Sediments.** The data (*SI Appendix, Fig. S2A*) on carbonate mineral sediment contents are composed of 1.7 million sites throughout the world's ocean, extracted from the comprehensive database of seabed sediment properties, dbSEABED (15). The original data are quality controlled and harmonized (16) so that analytical results and descriptions can be combined into a statistically homogeneous set of values. The sites are located geographically and by water depth. The surficial 0- to 10-cm (occasionally 0- to 30-cm) average values were used to compile the maps, a depth that takes account of scales of seafloor–seawater contact, including the effects of erosion and bioturbation.

**Water-Side  $\text{CO}_3^{2-}$  Mass Transfer Coefficient ( $\beta$ ).** The DBL thickness ( $Z_{\text{DBL}}$ ) at the air–sea interface is commonly derived from wind speed (42). In analogy, we describe a set of equations to estimate the benthic  $Z_{\text{DBL}}$  and the  $\text{CO}_3^{2-}$  water-side mass transfer coefficient ( $\beta$ ) at in situ  $T$  and  $S_p$ , using the horizontal current speed ( $U$ ) near the seafloor. In this study, we consider  $\beta$  to be independent of pressure, as the pressure effect on the viscosity of seawater and on the diffusion coefficients of ions in seawater is thought to be relatively small, that is, at most 8% (43). We use an interpolated version of horizontal bottom-water current speeds at  $1^\circ \times 1^\circ$  resolution from a higher resolution model (*SI Appendix, Fig.*

S3A). The bottom currents were inferred from an inline computation of the annually averaged kinetic energy field of a nominally 1/25th degree global configuration of the hybrid coordinate ocean model (HYCOM). It accounts for topographic internal lee wave drag, forced by air-sea fluxes and winds from reanalysis, but does not include tides (44). Neglecting tides is a source of uncertainty in this study, and this should be examined in future research. Nevertheless, abyssal tidal velocities (typically about  $1 \text{ cm s}^{-1}$ ) (45) are very likely to be smaller than nontidal velocities over most areas studied in this paper (SI Appendix, Fig. S3A). For a discussion of the sensitivity of our dissolution rate model to various initial kinetic energy distribution models, please see SI Appendix. Herein, for the sake of brevity, we refer to the current “speed” (a scalar) as current “velocity” (a vector), although our model does not use or require any information about the direction of the current. We assume (i) that these current velocities, averaged over the bottom 500 m, are representative of the current velocities just above the bottom boundary layer, that is, portions of sediment and water column affected directly in the distribution of their properties and processes by the presence of the SWI (13), comprising, among others, the DBL, (ii) that the seafloor is hydrodynamically smooth, and (iii) that there is no vertical density gradient, that is, the bottom boundary layer is unstratified.

Using the universal logarithmic velocity distribution, we then have the following:

$$\frac{u(z)}{u^*} = \frac{1}{\kappa} \ln \left( \frac{zu^*}{\nu} \right) + C, \quad [4]$$

where  $z$  is the height above the bottom,  $\kappa$  is the dimensionless von Karman constant (0.40) from ref. 46,  $\nu$  is the molecular kinematic viscosity of seawater at  $S_p = 35$  and in situ  $T$ , and  $C$  is an empirical dimensionless constant with a value of 5.1, taken from ref. 46.

The current velocity  $U$  at the upper edge of the bottom boundary layer is assumed to be equal to the current velocity averaged over the bottom 500 m and is then given by the following:

$$U = \frac{u^*}{\kappa} \ln \left( \frac{\delta u^*}{\nu} \right) + Cu^*, \quad [5]$$

where  $\delta$  is the bottom boundary layer thickness.

We further assume that  $\delta$  corresponds to the so-called Ekman scale height (13), reflecting classical Ekman dynamics, in which the height of the bottom boundary layer is limited by the Earth’s rotation:

$$\delta = \frac{\kappa u^*}{f}, \quad [6]$$

where  $f$  is the Coriolis parameter (set to  $1 \times 10^{-4} \text{ s}^{-1}$ ). We consider a constant  $S_p$  of 35, due to the small variations of this variable in oceanic bottom waters, as shown in SI Appendix, Fig. S1C.

Combining Eqs. 5 and 6, we obtain the following:

$$U = \frac{u^*}{\kappa} \ln \left( \frac{\kappa u^{*2}}{f\nu} \right) + Cu^*. \quad [7]$$

Fitting a power function that relates  $u^*$  to  $U$ , for  $T = 2^\circ \text{C}$ , with  $R^2 > 0.999$ , we have that

$$u^* = 0.023U^{0.84}. \quad [8]$$

This equation gives  $u^*$  values that are notably similar to previously published empirical relationships (47, 48), for example,  $u^* = U/30$ , as shown in SI Appendix, Fig. S4A.

The mass transfer coefficient,  $\beta$ , can be found for the carbonate ion ( $\text{CO}_3^{2-}$ ) following ref. 49 and  $u^*$  from Eq. 8:

$$\beta = 0.0417u^*Sc^{-2/3}, \quad [9]$$

where  $Sc$  is the Schmidt coefficient for  $\text{CO}_3^{2-}$ , that is,  $Sc = \nu/D_{\text{CO}_3^{2-}}$  at  $S_p = 35$  and in situ  $T$ , and  $D_{\text{CO}_3^{2-}}$  is the  $\text{CO}_3^{2-}$  diffusion coefficient from ref. 50. The geographical distribution of  $\beta$  is provided in SI Appendix, Fig. S3C. It should be noted that  $\beta$  is dependent on the temperature (SI Appendix, Fig. S4C), and, in contrast to the deep sea where the temperature is relatively constant (SI Appendix, Fig. S1B), significant variations can be encountered in warmer environments such as on carbonate platforms. Although studies such as ref. 51 have shown that there might have been a measurable change in bottom-water temperatures over the hydrographic period (late 1970s to today), accounting for it would have a very small impact on our results, given that these changes are currently limited to a few tenths of a degree Celsius (51).

Finally, we provide a  $1^\circ \times 1^\circ$  global model of the DBL thicknesses ( $Z_{\text{DBL}}$ ) at the SWI, obtained by dividing  $D_{\text{CO}_3^{2-}}$  (at  $S_p = 35$  and in situ  $T$ ) by  $\beta$ , in SI Appendix, Fig. S3B. This model can be used to estimate solute exchange at the SWI and is available on the NOAA Ocean Carbon Data System website, <https://www.nodc.noaa.gov/ocads/oceans/>.

Note that other empirical laws estimating  $\beta$  as function of  $Sc$  and  $u^*$ , similar to Eq. 9, can be found in the literature (52–55). In SI Appendix, Fig. S4B, we provide a comparison of the DBL thickness as a function of the shear velocity ( $u^*$ ) for  $T = 2^\circ \text{C}$  and  $S_p = 35$  from various studies. SI Appendix, Fig. S4 D–F shows the range of bottom-current velocities, shear velocities, DBL thicknesses, and  $\text{CO}_3^{2-}$  mass transfer coefficients encountered at the seafloor. In addition, we represent in SI Appendix, Fig. S4B the DBL thicknesses observed in ref. 19 based on the dissolution of alabaster plates deployed at the seafloor, as a function of the shear velocities from current-meter measurements or skin friction probes (19).

**Sediment-Side Mass Transfer Coefficient ( $k_s$ ).** Values for  $k_s$  have been obtained from experiments on the dissolution of  $\text{CaCO}_3$  beds, as reported in refs. 12 and 56. Based on ref. 12, whose results are statistically similar to those from ref. 56,

$$k_s \approx (326.9 \pm 39.5 \text{ m a}^{-1}) X_{\text{calcite}}^{0.5}, \quad [10]$$

at  $25^\circ \text{C}$ , where  $X_{\text{calcite}}$  is the calcite content in surficial sediments from the  $1^\circ \times 1^\circ$  distribution model in SI Appendix, Fig. S2A. The sediment-side mass transfer coefficients were subsequently adjusted from the laboratory temperature to the deep-seafloor temperature ( $T = 2^\circ \text{C}$ ) by assuming maximum activation energy for calcite dissolution in seawater of  $50 \text{ kJ mol}^{-1}$ , following the arguments presented in ref. 11. The distribution of  $k_s$  is shown in SI Appendix, Fig. S2B.

We provide a comparison of in situ  $k_s$  and  $\beta$  in SI Appendix, Fig. S5, where the blue end of the color bar represents a strongly water-side transport-controlled calcite dissolution reaction, while the red end stands for a strongly sediment-side controlled reaction. As it can be seen in SI Appendix, Fig. S5, calcite dissolution over most (72%) of the seafloor is water-side transport-controlled, except for areas covered by calcite-depleted sediments, where the reaction rate is limited by internal processes such as porewater diffusion or processes occurring at the surface of the calcite grains, thus independent of bottom-water hydrodynamics (see the discussion in SI Appendix for further details).

**CCD and Calcite Downward Flux at the CCD ( $F$ ).** The calcite content data in surficial sediments ( $X_{\text{calcite}}$ ) from the dbSEABED database were used to estimate the CCD in each basin, as defined in SI Appendix, Fig. S6, excluding the Arctic and Southern Oceans (below  $60^\circ \text{S}$ ), as well as continental shelves and near-shore areas. As mentioned in the Introduction, the CCD is operationally defined as the depth where the  $\text{CaCO}_3$  content falls below 10%. Thus, for each basin, we selected the data points for which  $X_{\text{calcite}}$  is equal to  $0.10 \pm 0.03$  (or  $10 \pm 3\%$ ) and computed the arithmetic average of the seafloor depths associated with these measurements. These depths constitute our CCD estimates for each basin and are assumed to be representative of preindustrial conditions. The CCD for each basin is reported along with the associated uncertainty and number of measurement used in SI Appendix, Table S1.

At the CCD, by definition, the calcite sinking rate is equal to its dissolution rate. Thus, the calcite downward flux at the CCD level ( $F$ ) was computed from the estimated preindustrial CCD and GLODAPv2.2016b data, per the following:

$$F = \frac{K_{sp}^* k^*}{[\text{Ca}^{2+}]} \exp \left( \text{CCD} \frac{\rho g}{p_c} \right) - k^* [\text{CO}_3^{2-}]_{\text{SW}}^{p-i}, \quad [11]$$

where  $p_c$  is a characteristic pressure (2) set to 511 atm,  $\rho$  is the mean density of seawater,  $g$  is the mean gravitational acceleration on Earth, and  $[\text{CO}_3^{2-}]_{\text{SW}}^{p-i}$  is the preindustrial bottom-water  $\text{CO}_3^{2-}$  concentration.

**Calcite Dissolution Rate Calculation.** For preindustrial conditions, the dissolution rate at the seafloor between the CSD and the CCD is given by the following:

$$r = k^* \left( [\text{CO}_3^{2-}]_{\text{eq}} - [\text{CO}_3^{2-}]_{\text{SW}} \right), \quad [12]$$

which is Eq. 2 of the text. For the seafloor located below the CCD, all  $\text{CaCO}_3$  arriving at the seafloor ( $F$ ) dissolves, that is,

$$r = F, \quad [13]$$

and these two latter equations account for all of the benthic dissolution.

With added anthropogenic dissolution, Eq. 12 still accounts for dissolution between the CSD and the CCD, both of which may now have been migrated

upward, and Eq. 13 still accounts for dissolution below the snowline, which remained in place. To account for dissolution between the snowline and the CCD, which have now separated as a result of anthropogenic CO<sub>2</sub> intrusion in the deep ocean, all CaCO<sub>3</sub> settling between these depths now dissolves at the seafloor (and in the water column). In addition, previously deposited CaCO<sub>3</sub> will dissolve as well. We account for these two sources by setting

$$r = F + k^* \left( [\text{CO}_3^{2-}]_{\text{eq}} - [\text{CO}_3^{2-}]_{\text{sw}} \right). \quad [14]$$

This assumes that dissolution has not been sufficient to deplete CaCO<sub>3</sub> in the sediment.

**Statistical Treatment.** SDs for GLODAPv2 variables (shown in *SI Appendix, Fig. S1*) were computed as the sum of the observational errors reported in ref. 57 and the mapping interpolation errors from ref. 26. The SD on the average CCD values are reported along with the number of measurements used for each estimate. SDs for  $[\text{CO}_3^{2-}]_{\text{sw}}$ ,  $\Delta[\text{CO}_3^{2-}]_{\text{sw}}$ ,  $r$ ,  $\Delta r$ , and  $F$  were computed

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